A FREE BOUNDARY PROBLEM FOR AN ELLIPTIC SYSTEM

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ABSTRACT. We study solutions and the free boundary $\partial\{|{\bf u}| > 0\}$ of the sublinear system

$$
\Delta \mathbf{u} = \lambda_+(x)|\mathbf{u}^+|^{q-1}\mathbf{u}^+ - \lambda_-(x)|\mathbf{u}^-|^{q-1}\mathbf{u}^-,
$$

from a regularity point of view.

For $\lambda_{\pm}(x) > 0$ and Hölder, and $0 < q < 1$, we apply the epiperimetric inequality approach and show $C^{1,\beta}$ -regularity for the free boundary at asymptotically flat points.

CONTENTS

1. INTRODUCTION

1.1. Problem setting. In this paper we study the elliptic system

(1)
$$
\Delta \mathbf{u} = \lambda_{+}(x)|\mathbf{u}^{+}|^{q-1}\mathbf{u}^{+} - \lambda_{-}(x)|\mathbf{u}^{-}|^{q-1}\mathbf{u}^{-},
$$

where $\lambda_{\pm} > 0$ are Hölder regular, $\mathbf{u} = (u_1, \ldots, u_m)$, with $\mathbf{u} : B_1 \subset \mathbb{R}^n \longrightarrow \mathbb{R}^m$, $n \geq 2, m \geq 1$, and $\mathbf{u}^{\pm} = (u_1^{\pm}, \cdots, u_m^{\pm})$. Here $|\cdot|$ stands for the Euclidian norm, $B_1 = B_1(0)$ is the unit ball, and the equation is in the weak sense. Solutions of [\(1\)](#page-0-2) are the unique minimizers (up to the prescribed boundary values) of the energy

(2)
$$
J_0(\mathbf{u}) = \int_{B_1} \left(|\nabla \mathbf{u}|^2 + \frac{2}{1+q} \lambda_+(x) |\mathbf{u}^+|^{1+q} + \frac{2}{1+q} \lambda_-(x) |\mathbf{u}^-|^{1+q} \right) dx.
$$

We are interested in the regularity of both minimizers \bf{u} of [\(2\)](#page-0-3) and their free boundaries $\partial\{x : |u(x)| > 0\}$. Our departing point is a $W^{1,2}$ -solution to this

equation, regardless of the boundary data. Since for each $i = 1, \dots, m$, $\Delta u_i \in \mathcal{L}^{2/q}$ we will have $u_i \in W^{2,2/q}$ and a bootstrap argument will show that $u_i \in W^{2,p}$ for all $p < \infty$. The main question is about the higher regularity of the solution along with the regularity of the free boundary $\partial\{|{\bf u}|>0\}$.

For clarity of exposition, and for readers' convenience we shall carry out the analysis for the case $\lambda_+ = \lambda_- \equiv 1$. In Section [5](#page-17-0) we shall explain the obvious and necessary changes for the general case in [\(1\)](#page-0-2). We thus, in what follows, consider the equation

(3)
$$
\Delta \mathbf{u} = f(\mathbf{u}) := |\mathbf{u}|^{q-1} \mathbf{u}, \quad \text{in } B_1(0), \quad \text{where } q \in (0, 1),
$$

that are minimizers to

(4)
$$
J(\mathbf{u}) = \int_{B_1} \left(|\nabla \mathbf{u}|^2 + 2F(\mathbf{u}) \right) dx,
$$

where

$$
F(\mathbf{u}) = \frac{1}{1+q} |\mathbf{u}|^{1+q}.
$$

The case $q = 0$ was studied in [\[2\]](#page-21-1) and it has been shown that the set of "regular"^{[1](#page-1-1)} free boundary points is locally a $C^{1,\beta}$ surface. It is noteworthy that in the scalar case, when $m = 1$, one recovers also the two phase free boundary problem

(5)
$$
\Delta u = (u^+)^q - (u^-)^q,
$$

which was investigated in [\[4\]](#page-21-2), from a regularity point of view, with partial results. When solutions of [\(5\)](#page-1-2) are assumed to be non-negative, the optimal regularity $C^{\lfloor \kappa \rfloor, \kappa-\lfloor \kappa \rfloor}$ for the solution has been shown, where $\kappa = 2/(1-q)$, as well as the regularity of the free boundary close to almost flat points; see [\[1,](#page-21-3) [3,](#page-21-4) [7,](#page-21-5) [8\]](#page-21-6).

In this paper we study the behaviour of solutions as well as the free boundary close to asymptotically flat points, and obtain results along the lines of [\[2\]](#page-21-1).

1.2. Notations and Definitions. For clarity of exposition we shall introduce some notation and definitions here that are used frequently in the paper.

Throughout this paper, \mathbb{R}^n will be equipped with the Euclidean inner product $x \cdot y$ and the induced norm |x|, $B_r(x_0)$ will denote the open *n*-dimensional ball of center x_0 , radius r with the boundary $\partial B_r(x_0)$. In addition, $B_r = B_r(0)$ and $\partial B_r = \partial B_r(0)$. For a set A, $d(x, A)$ stands for the distance between x and A. In the text we use the *n*-dimensional Hausdorff measure \mathcal{H}^n . For a real number s, we denote the greatest integer below s by $|s|$, i.e. $s - 1 \leq |s| < s$.

Also, we will denote the derivative of function f by $f_{\mathbf{u}}$ and the derivative matrix of **u** by ∇ **u** = $[\partial_i u_j]_{1 \leq i \leq n, 1 \leq j \leq m}$ with the notation

$$
|\nabla \mathbf{u}|^2 = \sum_{i=1}^m |\nabla u_i|^2, \qquad \nabla \mathbf{u} : \nabla \mathbf{v} = \sum_{i=1}^m (\nabla u_i \cdot \nabla v_i),
$$

\n
$$
\nabla \mathbf{u} \cdot \boldsymbol{\xi} = \boldsymbol{\xi}^t \nabla \mathbf{u} = (\nabla u_1 \cdot \boldsymbol{\xi}, \cdots, \nabla u_m \cdot \boldsymbol{\xi}), \qquad \text{for all } \boldsymbol{\xi} \in \mathbb{R}^n.
$$

We denote by $\Gamma(\mathbf{u}) = \partial \{|\mathbf{u}| > 0\} \cap {\{|\nabla \mathbf{u}| = 0\}}$ the set of free boundary. Moreover, for $q \in (0, 1)$ we fix the following constants throughout the paper:

$$
\kappa = \frac{2}{1-q}, \qquad \alpha = (\kappa(\kappa - 1))^{-\kappa/2}.
$$

¹For the meaning of regular points, see Definition [1.3.](#page-2-2)

1.3. Main results. Let us first assume that a solution $\mathbf{u} = (u_1, \ldots, u_m)$ of [\(3\)](#page-1-3), is such that all components of **u** are positive. Then we have the following result.

Proposition 1.1 (Regularity near the one-phase free boundary points). Let **u** be a solution of the system [\(3\)](#page-1-3), and $u_i \geq 0$ in $B_r(x_0)$ for some i. Then there is a constant $c = c(n, q)$ such that

$$
u_i(x) \le c(u_i(x_0) + |x - x_0|^{\kappa}), \quad \forall x \in B_{r/2}(x_0).
$$

This theorem shows that if $\mathbf{u}(x_0) = 0$, then all derivatives of **u** of order less than κ at point x_0 vanish. However, we can not expect to obtain $C^{\lfloor \kappa \rfloor, \kappa - \lfloor \kappa \rfloor}$ -regularity in the general case, particularly when some components of u change signs. Indeed, the ODE $y'' = y^q$, with initial condition $y(0) = 0 \neq y'(0)$, has a solution whose third derivative is unbounded, $y''' = qy'y^{q-1}$.

In order to study the optimal decay of solutions near such points, we start with a definition of the subset $\Gamma^{s}(\mathbf{u})$ of the free boundary $\Gamma(\mathbf{u})$ as follows

 $\Gamma^{s}(\mathbf{u}) := \{z \in \Gamma(\mathbf{u}) : \text{there exists some } c > 0 \text{ and a vector function } \mathbf{P}_m \text{ that each }$

component is a polynomial of degree at most $m < s$, such

that for all
$$
r > 0
$$
 we have $\sup_{B_r(0)} |\mathbf{u}(x+z) - \mathbf{P}_m(x)| \leq cr^s$.

We will show that $\Gamma^{\kappa}(\mathbf{u})$ contains only points at that all derivatives of order less than κ are zero.

Theorem 1.2. Let **u** be a solution of the system [\(3\)](#page-1-3) with $u(z) = 0$. Consider $\ell = |\kappa|$ to be the greatest integer below κ , i.e. $\kappa - 1 \leq \ell < \kappa$, then $z \in \Gamma^{\kappa}(\mathbf{u})$ if and only if

$$
\sup_{B_r(z)} |\mathbf{u}| = o(r^{\ell}).
$$

To investigate the regularity of free boundary, we consider "asymptotically onephase-points" that is, a subset of $\Gamma(\mathbf{u})$ such that the blow-ups belong to

 $\mathbb{H} := \{x \mapsto \alpha \max(x \cdot \nu, 0)^{\kappa} \mathbf{e} : \nu \in \mathbb{R}^n \text{ and } \mathbf{e} \in \mathbb{R}^m \text{ are unit vectors}\}.$

Members of these class are κ -homogeneous global solutions of [\(3\)](#page-1-3). When the domain is a plane, i.e. $n = 2$, all κ -homogeneous global solutions are classified (see Proposition [6.1](#page-20-1) in Appendix). The members of $\mathbb H$ are called *half-plane* solutions.

Definition 1.3. We denote by $\mathcal{R}_{\mathbf{u}}$ the set of all regular free boundary points of $x_0 \in \Gamma(\mathbf{u})$, which has at least one blow-up limit of **u** at x_0 that belongs to \mathbb{H} .

This definition is well-defined according to the uniqueness of blow-up, as we will show later. (See Remark [4.4\)](#page-14-0). Our main result concerning the regularity of the free boundary is presented in the following theorem.

Theorem 1.4 (Regularity of the free boundary). The set of regular free boundary points $\mathcal{R}_{\mathbf{u}}$ is locally in B_1 a $C^{1,\beta}$ -manifold.

We choose the epiperimetric inequality approach to prove this result. Since the first application of this approach in [\[9\]](#page-21-7), it has been used in various articles (see for example [\[2\]](#page-21-1) for an application in a system or [\[10\]](#page-21-8) for a sublinear scalar equation case). This inequality, Theorem [3.1,](#page-6-1) with a monotonicity formula, Proposition [2.3,](#page-4-0) provides an estimate for the energy decay. Indeed, one can control the rate of convergence $\|\mathbf{u}(x_0 + \cdot) - \mathbf{h}\|_{\mathcal{L}^1(B_r)}$ in this approach, where h belongs to H. In

Theorem [4.7,](#page-16-1) we will show that when $\mathbf{h} \in \mathbb{H}$, the rate of convergence is $r^{n+\kappa+\beta}$ for some $\beta > 0$.

In order to keep the presentation simple, we consider the constant coefficient case [\(3\)](#page-1-3) and do all calculation first for that. In Section [5,](#page-17-0) the result is extended to the general form [\(1\)](#page-0-2).

2. Higher Regularity of Solutions

In this section we will study the regularity of solutions of [\(3\)](#page-1-3) and prove Proposition [1.1](#page-2-3) and Theorem [1.2.](#page-2-4)

Proof of Proposition [1.1.](#page-2-3) Let $\varphi(r) := \int_{\partial B_r(x_0)} u_i$ and observe that $u_i \geq 0$, in the statement of the theorem. Since $\Delta u_i = |\mathbf{u}|^{q-1} u_i \geq 0$, we know that $\varphi(r)$ is increasing and

(6)
$$
\varphi'(r) = \frac{r}{n} \int_{B_r(x_0)} |\mathbf{u}|^{q-1} u_i \leq \frac{r}{n} \Big(\int_{B_r(x_0)} u_i \Big)^q \Big(\int_{B_r(x_0)} \frac{u_i}{|\mathbf{u}|} \Big)^{1-q} \leq \frac{r}{n} \Big(\varphi(r) \Big)^q,
$$

where in the last inequality we have applied $\int_{B_r(x_0)} u_i \leq \varphi(r)$. Form [\(6\)](#page-3-1), we obtain

$$
\varphi(r)^{1-q} - \varphi(0)^{1-q} \le \frac{r^2}{2n},
$$

and hence

$$
\varphi(r) \le (u_i(x_0)^{1-q} + \frac{r^2}{2n})^{\kappa/2} \le C_q(u_i(x_0) + r^{\kappa}).
$$

On the other hand, u_i is a nonnegative subharmonic function and there is a constant C_n such that

$$
u_i(x) \le C_n \int_{\partial B_\rho(x_0)} u_i, \quad \text{for every } x \in B_{r/2}(x_0) \text{ and } \rho = 2|x - x_0|.
$$

Therefore, $u_i(x) \le c(n, q)(u_i(x_0) + \rho^{\kappa}) \le c(u_i(x_0) + |x - x_0|^{\kappa}).$

.

Corollary 2.1. Let **u** be a solution of the system [\(3\)](#page-1-3) and $x_0 \in \Gamma$. There exists a constant $c = c(n, q)$ such that if $u_i \geq 0$ in $B_r(x_0)$ for all i, then

$$
\sup_{B_r(x_0)}|\mathbf{u}|\leq cr^{\kappa}
$$

Now we are going to prove Theorem [1.2.](#page-2-4) The necessity of vanishing derivatives is deduced from the following proposition when $s = \kappa$. The proof follows the same line of reasoning as that of the proof of Proposition 2.1 in [\[4\]](#page-21-2).

Proposition 2.2. Let **u** be a solution of system [\(3\)](#page-1-3) and $z \in \Gamma^s(\mathbf{u})$ for $s \leq \kappa$, then all the derivatives of order $m < \frac{s-2}{q}$ at point z are zero.

This proposition in a special case will imply that if $z \in \Gamma^{\kappa}(\mathbf{u})$ then all derivatives of **u** at the point z up to order $m < \kappa$ exist and are equal to zero, i.e. we must have $\mathbf{P}_m(x) = 0$ in definition of $\Gamma^{\kappa}(\mathbf{u})$. Thus when we are looking for points with regularity $C^{\lfloor \kappa \rfloor, \kappa-\lfloor \kappa \rfloor}$ in the free boundary $\Gamma(\mathbf{u})$, we might find them among the points where all the derivatives below κ exist and are zero. Theorem [1.2](#page-2-4) shows that this is a sufficient condition for a free boundary point to belong to $\Gamma^{\kappa}(\mathbf{u})$. We divide the proof in two different cases, depending on whether κ is integer or not. Before that we need to show that the monotonicity formula (which is established by the third author in [\[9\]](#page-21-7) for the classical obstacle problem), holds in the present setting.

Proposition 2.3. Let **u** be a solution of [\(3\)](#page-1-3) in $B_{r_0}(x_0)$ and let

$$
W(\mathbf{u},x_0,r) = \frac{1}{r^{n+2\kappa-2}} \int_{B_r(x_0)} \left(|\nabla \mathbf{u}|^2 + 2F(\mathbf{u}) \right) dx - \frac{\kappa}{r^{n+2\kappa-1}} \int_{\partial B_r(x_0)} |\mathbf{u}|^2 d\mathcal{H}^{n-1}.
$$

(i) For $0 < r < r_0$, the energy function $W(\mathbf{u}, x_0, r)$ is non-decreasing.

(ii) The function $x \mapsto W(\mathbf{u}, x, 0+)$ is upper-semicontinuous.

Proof. For $\mathbf{u}_r(x) := \mathbf{u}(x_0+rx)/r^{\kappa}$ we can apply the relations $r\partial_r \mathbf{u}_r = \nabla \mathbf{u}_r \cdot x - \kappa \mathbf{u}_r$ and $W(\mathbf{u}_r, 0, s) = W(\mathbf{u}, x_0, rs)$ to write for $s > t > 0$,

$$
W(\mathbf{u},x_0,s)-W(\mathbf{u},x_0,t)=\int_t^s\int_{\partial B_1(0)}\frac{2}{r}|\nabla \mathbf{u}_r\cdot x-\kappa\mathbf{u}_r|^2d\mathcal{H}^{n-1}dr\geq 0.
$$

For (ii), if $W(\mathbf{u}, x_0, 0+) > -\infty$, for an arbitrary $\varepsilon > 0$, we may by monotonicity, part (i), choose r such that $W(\mathbf{u}, x_0, r) \leq W(\mathbf{u}, x_0, 0+) + \varepsilon/2$. For this fixed r, there is a δ -neighborhood of x_0 such that $W(\mathbf{u}, x, r) \leq W(\mathbf{u}, x_0, r) + \varepsilon/2$. Therefore,

$$
W(\mathbf{u},x,0+)\leq W(\mathbf{u},x,r)\leq W(\mathbf{u},x_0,r)+\varepsilon/2\leq W(\mathbf{u},x_0,0+)+\varepsilon.
$$

The case, $W(\mathbf{u}, x_0, 0+) = -\infty$ will be proved by a similar argument. We must show that for an arbitrary constant $M > 0$, $W(\mathbf{u}, x, +0) < -M$ in some neighborhood of x_0 . Here, choose $r > 0$ such that $W(\mathbf{u}, x_0, r) \le -2M$, and for this r take a δ-neighborhood of x_0 such that $W(\mathbf{u}, x, r) \leq W(\mathbf{u}, x_0, r) + M \leq -M$.

Proof of sufficiency part of Theorem [1.2.](#page-2-4) Case $\kappa \notin \mathbb{N}$: If the statement of the theorem fails, then there exists a sequence $r_j \to 0$ such that

$$
\sup_{B_r} |\mathbf{u}| \leq j r^{\kappa}, \quad \forall \ r \geq r_j, \qquad \sup_{B_{r_j}} |\mathbf{u}| = j r_j^{\kappa}.
$$

In particular the function $\tilde{\mathbf{u}}_j(x) = \frac{\mathbf{u}(r_j x)}{jr_j^{\kappa}}$ satisfies

$$
\sup_{x \in B_R} |\tilde{\mathbf{u}}_j(x)| \le R^{\kappa}, \qquad \text{for } 1 \le R \le \frac{1}{r_j},
$$

with equality for $R = 1$, along with

$$
\Delta \tilde{\mathbf{u}}_j = \frac{\Delta \mathbf{u}(r_j x)}{j r_j^{\kappa - 2}} = \frac{f(\tilde{\mathbf{u}}_j)}{j^{1-q}} \longrightarrow 0
$$
 locally uniformly.

From here we conclude that $\{\tilde{\mathbf{u}}_j\}$ is bounded in $C^{1,\alpha}_{\text{loc}}(\mathbb{R}^n)$ and that there is a convergent subsequence, tending to a harmonic function \mathbf{u}_0 with growth κ , i.e.

(7)
$$
\sup_{B_R} |\mathbf{u}_0| \le R^{\kappa}, \quad \text{ for all } R \ge 1, \quad \sup_{B_1} |\mathbf{u}_0| = 1, \quad \Delta \mathbf{u}_0 = 0,
$$

and

(8)
$$
\mathbf{u}_0(0) = |\nabla \mathbf{u}_0(0)| = \cdots = |D^{\ell} \mathbf{u}_0(0)| = 0.
$$

Obviously [\(7\)](#page-4-1)-[\(8\)](#page-4-2), along with the fact that $\kappa \notin \mathbb{N}$, violates Liouville's theorem and we have a contradiction in this case.

Case $\kappa \in \mathbb{N}$: Let $\mathbf{u}_r(x) = \frac{\mathbf{u}(rx)}{r^{\kappa}}$ and notice that $\Delta \mathbf{u}_r = f(\mathbf{u}_r)$. By elliptic theory (see Theorem 8.17, in [\[6\]](#page-21-9)) it suffices to show that $\int_{B_1} |\mathbf{u}_r|^{1+q} dx$ is bounded. Using monotonicity formula, Theorem [2.3,](#page-4-0) we have

$$
\frac{2}{1+q} \int_{B_1} |\mathbf{u}_r|^{1+q} \le \int_{B_1} 2F(\mathbf{u}_r) \le W(\mathbf{u}_r, 0, 1) - \int_{B_1} |\nabla \mathbf{u}_r|^2 + \kappa \int_{\partial B_1} |\mathbf{u}_r|^2
$$

\n
$$
\le W(\mathbf{u}, 0, r) - \int_{B_1} |\nabla \mathbf{u}_r - \nabla \mathbf{p}|^2 + \kappa \int_{\partial B_1} |\mathbf{u}_r - \mathbf{p}|^2
$$

\n(9)
$$
\le W(\mathbf{u}, 0, 1) + \kappa \int_{\partial B_1} |\mathbf{u}_r - \mathbf{p}|^2,
$$

where each component of $\mathbf{p} = (p_1, \ldots, p_m) \in \mathbb{P}_{\kappa}$ is an arbitrary homogeneous harmonic polynomial of order κ . We need only to show that $\int_{\partial B_1} |\mathbf{u}_r - \pi_r|^2$ is bounded for every $r \leq 1$, where $\pi_r = \operatorname{argmin}_{\mathbf{p} \in \mathbb{P}_\kappa} \int_{\partial B_1} |\mathbf{u}_r - \mathbf{p}|^2$. The function π_r satisfies Z

$$
\int_{\partial B_1} \mathbf{p} \cdot (\mathbf{u}_r - \pi_r) d\mathcal{H}^{n-1} = 0, \quad \text{for every } \mathbf{p} \in \mathbb{P}_{\kappa}.
$$

Now suppose, towards a contradiction, that there is a sequence $r_k \to 0$, such that

$$
M_k = \left(\int_{\partial B_1} |\mathbf{u}_{r_k} - \pi_{r_k}|^2 d\mathcal{H}^{n-1}\right)^{1/2} \longrightarrow \infty.
$$

For
$$
\mathbf{w}_k = \frac{\mathbf{u}_{r_k} - \pi_{r_k}}{M_k}
$$
, we have $\|\mathbf{w}_k\|_{\mathcal{L}^2(\partial B_1)} = 1$ and $\Delta \mathbf{w}_k = \frac{f(\mathbf{u}_{r_k})}{M_k}$,

$$
\int_{B_1} |\Delta \mathbf{w}_k|^{(1+q)/q} \le \frac{1}{M_k^{(1+q)/q}} \int_{B_1} |\mathbf{u}_{r_k}|^{1+q} \le \frac{C}{M_k^{(1+q)/q}} \left(1 + \int_{\partial B_1} |\mathbf{u}_{r_k} - \pi_{r_k}|^2\right) \to 0,
$$

where in the last inequality we have used [\(9\)](#page-5-0). Now $\{w_k\}$ being bounded in $W^{2,2}(B_1)$ there is a weakly convergence subsequence with limit \mathbf{w}_0 , satisfying $\Delta \mathbf{w}_0 = 0$, $\|\mathbf{w}_0\|_{\mathcal{L}^2(\partial B_1)} = 1$ and

(10)
$$
\int_{\partial B_1} \mathbf{p} \cdot \mathbf{w}_0 = 0, \quad \text{for every } \mathbf{p} \in \mathbb{P}_{\kappa}.
$$

On the other hand, we have

$$
\int_{B_1} |\nabla \mathbf{w}_k|^2 - \kappa \int_{\partial B_1} |\mathbf{w}_k|^2 = \frac{1}{M_k^2} \Big(\int_{B_1} |\nabla \mathbf{u}_{r_k}|^2 - \kappa \int_{\partial B_1} |\mathbf{u}_{r_k}|^2 \Big)
$$

$$
\leq \frac{1}{M_k^2} \Big(W(\mathbf{u}, 0, r_k) - 2 \int_{B_1} F(\mathbf{u}_{r_k}) dx \Big)
$$

$$
\leq \frac{1}{M_k^2} W(\mathbf{u}, 0, 1) \longrightarrow 0.
$$

Therefore, we obtain

(11)
$$
\int_{B_1} |\nabla \mathbf{w}_0|^2 - \kappa \int_{\partial B_1} |\mathbf{w}_0|^2 \le 0.
$$

On the other hand by Lemma 4.1 in [\[11\]](#page-21-10), each component w_0^i of w_0 must satisfy

$$
\kappa \int_{\partial B_1} (w_0^i)^2 \le \int_{B_1} |\nabla w_0^i|^2.
$$

Summing over i and comparing with (11) , this along with

$$
w_0^i(0) = |\nabla w_0^i(0)| = \dots = |D^{\ell}w_0^i(0)| = 0
$$

implies that w_0^i is a homogeneous harmonic polynomial of order κ . But [\(10\)](#page-5-2) implies that $\mathbf{w}_0 = 0$ on ∂B_1 which contradicts $\|\mathbf{w}_0\|_{\mathcal{L}^2(\partial B_1)} = 1$. **Remark 2.4.** The first part of the proof, case $\kappa \notin \mathbb{N}$, works when the equation is *relaxed to* $|\Delta \mathbf{u}| \leq c_0 |\mathbf{u}|^q$.

3. The Epiperimetric Inequality

This section is devoted to provide the main tool of our approach, the epiperimetric inequality. Firstly, let us define the boundary adjusted energy

$$
M(\mathbf{v}) := \int_{B_1} \left(|\nabla \mathbf{v}|^2 + 2F(\mathbf{v}) \right) dx - \kappa \int_{\partial B_1} |\mathbf{v}|^2 d\mathcal{H}^{n-1}.
$$

Theorem 3.1 (The epiperimetric inequlaity). There exist $\varepsilon \in (0,1)$ and $\delta > 0$ such that if $c \in W^{1,2}(B_1;\mathbb{R}^m)$ is a homogeneous function of degree κ and $\|c - c\|$ $\mathbf{h} \Vert_{W^{1,2}(B_1;\mathbb{R}^m)} \leq \delta$ for some $\mathbf{h} \in \mathbb{H}$, then there exists a function $\mathbf{v} \in W^{1,2}(B_1;\mathbb{R}^m)$ such that $\mathbf{v} = \mathbf{c}$ on ∂B_1 and $M(\mathbf{v}) \leq (1 - \varepsilon)M(\mathbf{c}) + \varepsilon M(\mathbf{h}).$

Proof. Suppose toward a contradiction that there are sequences $\varepsilon_i \to 0$, $\delta_i \to 0$, $\mathbf{c}_i \in W^{1,2}(B_1;\mathbb{R}^m)$ and $\mathbf{h}_i \in \mathbb{H}$ such that \mathbf{c}_i is homogeneous of degree κ and satisfies

$$
\|\mathbf{c}_i - \mathbf{h}_i\|_{W^{1,2}(B_1; \mathbb{R}^m)} = \inf_{\mathbf{h} \in \mathbb{H}} \|\mathbf{c}_i - \mathbf{h}\| = \delta_i,
$$

and

(12)
$$
M(\mathbf{v}) > (1 - \varepsilon_i)M(\mathbf{c}_i) + \varepsilon_i M(\mathbf{h}_i)
$$
, for all $\mathbf{v} \in \mathbf{c}_i + W_0^{1,2}(B_1; \mathbb{R}^m)$.

Rotating in \mathbb{R}^n and in \mathbb{R}^m if necessary, we may assume that

$$
\mathbf{h}_i(x) = \alpha (x_n^+)^{\kappa} \mathbf{e}_1 =: \mathbf{h}(x),
$$

where $\mathbf{e}_1 = (1, 0, \dots, 0) \in \mathbb{R}^m$. Notice that the energy M takes a constant value on \mathbb{H} , and that subtracting $M(\mathbf{h})$ from the inequality [\(12\)](#page-6-2), we obtain

(13) $(1 - \varepsilon_i)(M(\mathbf{c}_i) - M(\mathbf{h})) < M(\mathbf{v}) - M(\mathbf{h}), \quad \text{for all } \mathbf{v} \in \mathbf{c}_i + W_0^{1,2}(B_1; \mathbb{R}^m).$ Now observe that for all $\phi = (\phi_1, \dots, \phi_m) \in W^{1,2}(B_1; \mathbb{R}^m)$,

$$
\delta M(\mathbf{h})(\phi) := 2 \int_{B_1} \nabla \mathbf{h} : \nabla \phi + |\mathbf{h}|^{q-1} \mathbf{h} \cdot \phi \, dx - 2\kappa \int_{\partial B_1} \mathbf{h} \cdot \phi \, d\mathfrak{H}^{n-1}
$$

=
$$
2 \int_{B_1} \left(-\Delta \mathbf{h} + f(\mathbf{h}) \right) \cdot \phi \, dx + 2 \int_{\partial B_1} \left(\nabla \mathbf{h} \cdot x - \kappa \mathbf{h} \right) \cdot \phi \, d\mathfrak{H}^{n-1} = 0.
$$

Thus we can subtract $(1-\varepsilon_i)\delta M(\mathbf{h})(\mathbf{c}_i-\mathbf{h})$ from the left hand side and $\delta M(\mathbf{h})(\mathbf{v}-\mathbf{h})$ from the right hand side of [\(13\)](#page-6-3) to obtain

(14)
$$
(1 - \varepsilon_i) \Big(\int_{B_1} |\nabla(\mathbf{c}_i - \mathbf{h})|^2 dx - \kappa \int_{\partial B_1} |\mathbf{c}_i - \mathbf{h}|^2 d\mathcal{H}^{n-1} + 2 \int_{B_1} F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) dx \Big) < \int_{B_1} |\nabla(\mathbf{v} - \mathbf{h})|^2 dx - \kappa \int_{\partial B_1} |\mathbf{v} - \mathbf{h}|^2 d\mathcal{H}^{n-1} + 2 \int_{B_1} F(\mathbf{v}) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{v} - \mathbf{h}) dx.
$$

Define now the normalized functions $\mathbf{w}_i := (\mathbf{c}_i - \mathbf{h})/\delta_i$, which along a subsequence converge weakly in $W^{1,2}(B_1;\mathbb{R}^m)$ to a function w. The proof proceeds then in the following four steps:

Step 1. $\mathbf{w} = 0$ in $B_1 \cap \{x_n < 0\}.$ Step 2. w solves the equation $\Delta \mathbf{w} = f_{\mathbf{u}}(\mathbf{h})(\mathbf{w})$ in $B_1 \cap \{x_n > 0\}$. Step 3. $\mathbf{w} \equiv 0$. Step 4. $\mathbf{w}_i \to 0$ strongly in $W^{1,2}(B_1; \mathbb{R}^m)$ as the subsequence $i \to \infty$.

Since $\|\mathbf{w}_i\|_{W^{1,2}(B_1;\mathbb{R}^m)} = 1$, Step 3 and Step 4 imply a contradiction proving the theorem.

Step 1. We insert $\mathbf{v} := (1 - \eta)\mathbf{c}_i + \eta \mathbf{h}$ in [\(14\)](#page-6-4) where $\eta \in W_0^{1,2}(B_1)$ is radially symmetric and satisfies $0 \leq \eta \leq 1$, and obtain

$$
(1 - \varepsilon_i) \int_{B_1} F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) dx
$$

$$
< C\delta_i^2 + \int_{B_1} F(\mathbf{v}) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{v} - \mathbf{h}) dx
$$

$$
\le C\delta_i^2 + \int_{B_1} (1 - \eta) (F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h})) dx,
$$

where convexity of F is used in the last inequality. From this we obtain

(15)
$$
\int_{B_1} (\eta - \varepsilon_i) \big(F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) \big) dx < C \delta_i^2.
$$

We first integrate on $B_1^+ = B_1 \cap \{x_n > 0\}$ where $|\mathbf{h}| > 0$, to arrive at

$$
\int_{B_1^+} (\eta - \varepsilon_i) (F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h})) dx
$$
\n
$$
= \int_{B_1^+} \int_0^1 \int_0^t (\eta - \varepsilon_i) (f_\mathbf{u}(\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})) (\mathbf{c}_i - \mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h})) ds dt dx
$$
\n
$$
= \int_{B_1^+} \int_0^1 \int_0^t (\eta - \varepsilon_i) (|\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{c}_i - \mathbf{h}|^2
$$
\n
$$
+ (q - 1)|\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-3} ((\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})) \cdot (\mathbf{c}_i - \mathbf{h}))^2) ds dt dx
$$
\n
$$
= \Big(\int_0^1 (\eta(r) - \varepsilon_i) r^{k(q+1)+n-1} dr \Big) \int_{\partial B_1^+} \int_0^1 \int_0^t (|\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{c}_i - \mathbf{h}|^2
$$
\n
$$
+ (q - 1)|\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-3} ((\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})) \cdot (\mathbf{c}_i - \mathbf{h}))^2) ds dt d\mathcal{H}^{n-1}
$$
\n
$$
\geq \Big(\int_0^1 (\eta(r) - \varepsilon_i) r^{k(q+1)+n-1} dr \Big) \int_{\partial B_1^+} \int_0^1 \int_0^t q |\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{c}_i - \mathbf{h}|^2 ds dt d\mathcal{H}^{n-1}
$$
\n
$$
= \int_{B_1^+} \int_0^1 \int_0^t q(\eta - \varepsilon_i) |\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{c}_i
$$

It is noteworthy that the phrase inside the parentheses in the last inequality is positive when ε_i is small enough. Now comparing the last inequality with [\(15\)](#page-7-0) gives us the bound

$$
\int_{B_1^-} \frac{(\eta-\varepsilon_i)}{\delta_i^{1-q}} F(\mathbf{w}_i) dx + \int_{B_1^+} \int_0^1 \int_0^t q(\eta-\varepsilon_i) |\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{w}_i|^2 ds dt dx < C.
$$

As $i\to\infty$ we conclude that

$$
\int_{B_1^-} \eta F(\mathbf{w}) dx \le 0.
$$

Therefore $\mathbf{w} = 0$ in B_1^- .

Step 2. Insert $\mathbf{v} := \eta(\mathbf{h} + \delta_i \mathbf{g}) + (1 - \eta)\mathbf{c}_i$ into [\(14\)](#page-6-4) where $\eta \in C_0^{\infty}(B_1^+)$ with values in [0, 1] and $\mathbf{g} \in \mathbf{W}^{1,2}(B_1; \mathbb{R}^m)$

$$
\int_{B_1} |\nabla \mathbf{w}_i|^2 dx + \frac{2}{\delta_i^2} \int_{B_1} F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) dx
$$

$$
< C\varepsilon_i + \int_{B_1} |\nabla ((1 - \eta)\mathbf{w}_i + \eta \mathbf{g})|^2 dx
$$

$$
+ \frac{2}{\delta_i^2} \int_{B_1} F(\mathbf{v}) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{v} - \mathbf{h}) dx.
$$

It follows that

$$
\int_{B_1} (1 - (1 - \eta)^2) |\nabla \mathbf{w}_i|^2 dx + \frac{2}{\delta_i^2} \int_{B_1^+} F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) dx
$$

<
$$
< C\varepsilon_i + \int_{B_1} |\nabla(\eta \mathbf{g})|^2 + 2\nabla((1 - \eta)\mathbf{w}_i) \cdot \nabla(\eta \mathbf{g})
$$

$$
+ |\nabla \eta|^2 |\mathbf{w}_i|^2 - 2(1 - \eta)(\nabla \mathbf{w}_i \cdot \nabla \eta) \cdot \mathbf{w}_i dx
$$

$$
+ \frac{2}{\delta_i^2} \int_{B_1^+} F(\mathbf{v}) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{v} - \mathbf{h}) dx,
$$

and then passing to the limit as $i \to \infty$,

$$
\int_{B_1} |\nabla \mathbf{w}|^2 dx + \limsup_{i \to \infty} \int_{B_1^+} \int_0^1 \int_0^t 2f_\mathbf{u}(\mathbf{h} + s\delta_i \mathbf{w}_i)(\mathbf{w}_i) \cdot \mathbf{w}_i ds dt dx
$$

\n
$$
\leq \int_{B_1} |\nabla ((1 - \eta)\mathbf{w} + \eta \mathbf{g})|^2 dx
$$

\n
$$
+ \liminf_{i \to \infty} \int_{B_1^+} \int_0^1 \int_0^t 2f_\mathbf{u}(\mathbf{h} + s\delta_i((1 - \eta)\mathbf{w}_i + \eta \mathbf{g}))
$$

\n(16)
$$
((1 - \eta)\mathbf{w}_i + \eta \mathbf{g}) \cdot ((1 - \eta)\mathbf{w}_i + \eta \mathbf{g}) ds dt dx.
$$

On the other hand, in $B_1^+ \cap \text{supp } \eta$, we have

$$
f_{\mathbf{u}}(\mathbf{h} + s\delta_i \mathbf{w}_i)(\mathbf{w}_i) \cdot \mathbf{w}_i = |\mathbf{h} + s\delta_i \mathbf{w}_i|^{q-1} |\mathbf{w}_i|^2
$$

+ $(q-1)|\mathbf{h} + s\delta_i \mathbf{w}_i|^{q-3} ((\mathbf{h} + s\delta_i \mathbf{w}_i) \cdot \mathbf{w}_i)^2$
 $\longrightarrow f_{\mathbf{u}}(\mathbf{h})(\mathbf{w}) \cdot \mathbf{w},$

where the convergence is valid due to the dominated convergence theorem. A similar convergence holds for the right hand side of [\(16\)](#page-8-0), and hence

$$
\int_{B_1} |\nabla \mathbf{w}|^2 dx + \int_{B_1^+ \cap \operatorname{supp} \eta} f_{\mathbf{u}}(\mathbf{h})(\mathbf{w}) \cdot \mathbf{w} dx
$$
\n
$$
\leq \int_{B_1} |\nabla ((1 - \eta)\mathbf{w} + \eta \mathbf{g})|^2 dx
$$
\n
$$
+ \int_{B_1^+ \cap \operatorname{supp} \eta} f_{\mathbf{u}}(\mathbf{h}) ((1 - \eta)\mathbf{w} + \eta \mathbf{g}) \cdot ((1 - \eta)\mathbf{w} + \eta \mathbf{g}) dx.
$$

Consider an open ball $B \subset B_1 \cap \{x_n > 0\}$. We may choose $\eta := 1$ in B and $\mathbf{g} := \mathbf{w}$ outside B to obtain that

$$
\int_B |\nabla \mathbf{w}|^2 dx + \int_B f_\mathbf{u}(\mathbf{h})(\mathbf{w}) \cdot \mathbf{w} dx \le \int_B |\nabla \mathbf{g}|^2 dx + \int_B f_\mathbf{u}(\mathbf{h})(\mathbf{g}) \cdot \mathbf{g} dx,
$$

for all $\mathbf{g} \in W^{1,2}(B; \mathbb{R}^m)$ coinciding with w on ∂B . Therefore,

 $\Delta \mathbf{w} = |\mathbf{h}|^{q-1} \mathbf{w} + (q-1)|\mathbf{h}|^{q-3} (\mathbf{h} \cdot \mathbf{w}) \mathbf{h}.$

Step 3. Let $w_j := \mathbf{w} \cdot \mathbf{e}_j$ for $1 \leq j \leq m$, then

$$
\Delta w_j = q\kappa(\kappa - 1)(x_n^+)^{-2}w_j
$$
, for $j = 1$,

and

$$
\Delta w_j = \kappa (\kappa - 1)(x_n^+)^{-2} w_j
$$
, for $j > 1$.

Now extend w_j to a homogeneous function of degree κ in $\{x_n > 0\}$ and define

$$
\tilde{w}_j(x',x_n) := \begin{cases} w_j(x',x_n), & x_n > 0, \\ -w_j(x',-x_n), & x_n < 0, \end{cases}
$$

which is a homogeneous weak solution of degree κ and satisfies

(17)
$$
\Delta \tilde{w}_j = \begin{cases} q\kappa(\kappa - 1)|x_n|^{-2}\tilde{w}_j, & \text{for } j = 1, \\ \kappa(\kappa - 1)|x_n|^{-2}\tilde{w}_j, & \text{for } j > 1. \end{cases}
$$

Note that as a result of Step 1, the trace of **w** vanishes on $\{x_n = 0\}$. If we consider now any multiindex $\mu \in \mathbb{Z}_+^{n-1} \times \{0\}$ and the higher order partial derivatives $\partial^{\mu}\tilde{w}_j = \zeta$ then ζ satisfies again in the same equation in \mathbb{R}^n . Then ζ is by repeated local energy estimates contained in $W_{loc}^{1,2}(\mathbb{R}^n)$ and ζ is a homogeneous function of degree $\kappa - |\mu|_1$. From the integrability and homogeneity we infer that $\partial^{\mu} \tilde{w}_j \equiv 0$ for $\kappa - |\mu|_1 - 1 \leq -n/2$. Thus $x' \mapsto \tilde{w}_j(x', x_n)$ is a polynomial and the homogeneity and integrability imply the existence of a polynomial p of degree $\deg p < \kappa + \frac{1}{2} - 1$, such that $w_j(x', x_n) = x_n^{\kappa} p(\frac{x'}{x_n})$ $\frac{x'}{x_n}$ for $x_n > 0$. Next we take $\mu \in \mathbb{Z}_+^{n-1} \times \{0\}$ such that $|\mu|_1 = \deg p$, then $\partial^{\mu} w_j = \partial^{\mu} p x_n^{\kappa-|\mu|_1}$. Comparing with the equation [\(17\)](#page-9-0), in the case that $\partial^{\mu} p \neq 0$, implies that

$$
(\kappa - |\mu|_1)(\kappa - |\mu|_1 - 1) = q\kappa(\kappa - 1), \text{ for } j = 1,
$$

$$
(\kappa - |\mu|_1)(\kappa - |\mu|_1 - 1) = \kappa(\kappa - 1), \text{ for } j > 1,
$$

and hence

$$
|\mu|_1 = 1
$$
, or $2\kappa - 2$, for $j = 1$,
 $|\mu|_1 = 0$, or $2\kappa - 1$, for $j > 1$.

On the other hand, we have deg $p < \kappa - \frac{1}{2}$, and only the cases $|\mu|_1 = 1$ for $j = 1$ and $|\mu|_1 = 0$ for $j > 1$ are possible. For $j = 1$, we obtain that $w_1(x) = x_n^{\kappa} (d + \ell \cdot x'/x_n)$, whereupon the equation for w_1 yields that

$$
\Delta w_1 = (\kappa - 1)(\kappa - 2)x_n^{\kappa - 3}(dx_n + \ell \cdot x') + 2d(\kappa - 1)x_n^{\kappa - 2} = q\kappa(\kappa - 1)x_n^{\kappa - 3}(dx_n + \ell \cdot x').
$$

We deduce that $d = 0$ and that $w_1(x) = x_n^{\kappa-1} \ell \cdot x'$ in $\{x_n > 0\}$. By similar argument for $j > 1$, we find that $\mathbf{w}(x) = (x_n^{\kappa-1}\ell_1 \cdot x', \ell_2 x_n^{\kappa}, \ldots, \ell_m x_n^{\kappa}).$

Recall that we have chosen **h** as the best approximation of \mathbf{c}_i in H. It follows that for $\mathbf{h}_{\nu}(x) := \alpha \max(x \cdot \nu, 0)^{\kappa} \mathbf{e}_1$,

(18)
$$
(\mathbf{w}_i, \mathbf{h}_{\nu} - \mathbf{h})_{W^{1,2}(B_1; \mathbb{R}^m)} \leq \frac{1}{2\delta_i} \|\mathbf{h}_{\nu} - \mathbf{h}\|_{W^{1,2}(B_1; \mathbb{R}^m)}^2.
$$

Now let $\nu \to \mathbf{e}_n$ so that $\frac{\nu-\mathbf{e}_n}{|\nu-\mathbf{e}_n|}$ converges to the vector ξ (where $\xi \cdot \mathbf{e}_n = 0$), then

$$
o(1) \ge \int_{B_1} (\mathbf{w}_i \cdot \mathbf{e}_1) \kappa(x_n^+)^{\kappa - 1} (x \cdot \xi) +
$$

$$
\nabla(\mathbf{w}_i \cdot \mathbf{e}_1) \cdot \left[\kappa(x_n^+)^{\kappa - 1} \xi + \kappa(\kappa - 1)(x_n^+)^{\kappa - 2} (x \cdot \xi) \mathbf{e}_n \right] dx.
$$

Choosing $\xi = (\ell_1, 0)$ and passing to the limit in i, we obtain that

$$
0 \ge \int_{B_1} \kappa (x_n^+)^{2\kappa - 2} (x' \cdot \ell_1)^2 + \kappa (x_n^+)^{2\kappa - 2} |\ell_1|^2
$$

+ $\kappa (\kappa - 1)^2 (x_n^+)^{2\kappa - 4} (x' \cdot \ell_1)^2 dx.$

Hence, $\ell_1 = 0$, and $\mathbf{w} \cdot \mathbf{e}_1 = 0$.

It remains to show that $\mathbf{w} \cdot \mathbf{e}_j = 0$ for $j > 1$. Apply once more the relation [\(18\)](#page-10-0) for $\mathbf{h}_t = \alpha(x_n^+)^{\kappa} \mathbf{e}_t$ instead of \mathbf{h}_v , where $\mathbf{e}_t = (\cos t) \mathbf{e}_1 \pm (\sin t) \mathbf{e}_j$, and let $t \to 0$. We obtain

$$
(\mathbf{w}_i, \pm \alpha (x_n^+)^{\kappa} \mathbf{e}_j)_{W^{1,2}(B_1; \mathbb{R}^m)} \leq 0.
$$

Therefore,

$$
\ell_j \|(x_n^+)^\kappa\|_{W^{1,2}(B_1)}^2 = 0,
$$

and $\ell_j = 0$.

Step 4. In order to show the strong convergence of \mathbf{w}_i in $W^{1,2}(B_1;\mathbb{R}^m)$, choose $\mathbf{v} := (1 - \eta)\mathbf{c}_i + \eta \mathbf{h}$ as a test function in [\(14\)](#page-6-4), where $\eta = \max(0, \min(1, 2(1 - |x|))).$ Then as in Step 1, we obtain

$$
\int_{B_1} |\nabla \mathbf{w}_i|^2 dx + \int_{B_1} \frac{\eta - \varepsilon_i}{\delta_i^2} \left[F(\mathbf{c}_i) - F(\mathbf{h}) - f(\mathbf{h}) \cdot (\mathbf{c}_i - \mathbf{h}) \right] dx
$$

\n
$$
\leq C \varepsilon_i + \int_{B_1} |\nabla ((1 - \eta) \mathbf{w}_i)|^2 dx
$$

\n
$$
= C \varepsilon_i + \int_{B_1} (1 - \eta)^2 |\nabla \mathbf{w}_i|^2 - 2(1 - \eta) (\nabla \mathbf{w}_i \cdot \nabla \eta) \cdot \mathbf{w}_i + |\nabla \eta|^2 |\mathbf{w}_i|^2 dx
$$

and

$$
\int_{B_{1/2}} |\nabla \mathbf{w}_i|^2 dx + \int_{B_1^-} \eta F(\mathbf{w}_i) dx + \int_{B_1^+} \int_0^1 \int_0^t q\eta |\mathbf{h} + s(\mathbf{c}_i - \mathbf{h})|^{q-1} |\mathbf{w}_i|^2 ds dt dx
$$

$$
\leq C\varepsilon_i + \int_{B_1} |\nabla \eta|^2 |\mathbf{w}_i|^2 - 2(1 - \eta)(\nabla \mathbf{w}_i \cdot \nabla \eta) \cdot \mathbf{w}_i dx.
$$

At this point, notice that second and third integral in the last relation are positive and use the homogeneity of w_i to obtain

$$
\int_{B_1} |\nabla \mathbf{w}_i|^2 dx = 2^{n+2\kappa-2} \int_{B_{1/2}} |\nabla \mathbf{w}_i|^2 dx
$$
\n
$$
\leq 2^{n+2\kappa-2} \Big(C\varepsilon_i + \int_{B_1} |\nabla \eta|^2 |\mathbf{w}_i|^2 - 2(1-\eta)(\nabla \mathbf{w}_i \cdot \nabla \eta) \cdot \mathbf{w}_i dx \Big) \to 0.
$$

4. Regularity of Free Boundary at Regular Points

In this section, we will study the regularity of the free boundary near a one-phase point at which at least one blow-up limit coincides with a half-plane solution.

4.1. Nondegeneracy.

Proposition 4.1. Let **u** be a solution of [\(3\)](#page-1-3) with $0 < q < 1$. Then there is a positive constant $c = c(q, n)$ such that if $x_0 \in \overline{\{|u| > 0\}}$ and $B_r(x_0) \subset B_1$, then

$$
\sup_{B_r(x_0)} |\mathbf{u}| \geq c r^{\kappa}
$$

.

Proof. Let $v(x) := |u(x)|^{1-q}$. Then

$$
\Delta v = (1 - q) + (1 - q) \frac{|\nabla \mathbf{u}|^2}{v^{\kappa - 1}} - \frac{1 + q}{1 - q} \frac{|\nabla v|^2}{v}.
$$

For any $y \in \{|\mathbf{u}| > 0\}$ (close to x_0), set $w(x) = c|x - y|^2$ for small constant $c > 0$ to be specified later. Then $h = v - w$ satisfies in $\{|u| > 0\}$

$$
\mathcal{L}h := \Delta h + \frac{1+q}{1-q} \left(\frac{\nabla(v+w)}{v} \cdot \nabla h - \frac{4c}{v} h \right) = (1-q) - 4c \left(\frac{n}{2} + \frac{1+q}{1-q} \right) + (1-q) \frac{|\nabla \mathbf{u}|^2}{v^{\kappa - 1}} \ge 0,
$$

provided that c is small enough. In particular h cannot attain a local maximum in $B_r(y) \cap {\{\mathbf{u}\}\geq 0\}}.$ On the other hand $h < 0$ on $\partial{\{\mathbf{u}\}\geq 0\}}$ and hence the positive maximum of h is attained on $\partial B_r(y)$, and we conclude that

$$
\sup_{\partial B_r(y)\cap\{|{\mathbf u}|>0\}}(v-w)\geq v(y)>0,
$$

which amounts to

$$
\sup_{\partial B_r(y) \cap \{|{\bf u}|>0\}} v \geq cr^2.
$$

Letting $y \to x_0$, we arrive at the statement of the lemma.

4.2. Energy decay.

Theorem 4.2 (Energy decay). Let $x_0 \in B_1 \cap \partial \{|u| > 0\}$, and suppose that the epiperimetric inequality holds with $\varepsilon \in (0,1)$ for each

$$
\mathbf{c}_r(x) := |x|^{\kappa} \mathbf{u}_r(\frac{x}{|x|}) = \frac{|x|^{\kappa}}{r^{\kappa}} \mathbf{u}(x_0 + \frac{r}{|x|}x)
$$

and for all $r \le r_0 < 1$. Finally let \mathbf{u}_0 denote an arbitrary blow-up limit of \mathbf{u} at x_0 and $\Lambda = (n + 2\kappa - 2)\varepsilon/(1 - \varepsilon)$. Then

$$
|W(\mathbf{u},x_0,r)-W(\mathbf{u},x_0,0+)|\leq |W(\mathbf{u},x_0,r_0)-W(\mathbf{u},x_0,0+)|\left(\frac{r}{r_0}\right)^{\Lambda},
$$

for $r \in (0, r_0)$, and there exists a constant C depending only on n and ε such that

$$
\int_{\partial B_1} |\mathbf{u}_r(x) - \mathbf{u}_0(x)| d\mathcal{H}^{n-1} \le C|W(\mathbf{u}, x_0, r_0) - W(\mathbf{u}, x_0, 0+)|^{1/2} \left(\frac{r}{r_0}\right)^{\Lambda/2}.
$$

Proof. We define

$$
e(r) := W(\mathbf{u}, x_0, r) - W(\mathbf{u}, x_0, 0+) = r^{-n-2\kappa+2} \int_{B_r(x_0)} \left(|\nabla \mathbf{u}|^2 + 2F(\mathbf{u}) \right) dx
$$

$$
-\kappa r^{-n-2\kappa+1} \int_{\partial B_r(x_0)} |\mathbf{u}|^2 d\mathcal{H}^{n-1} - W(\mathbf{u}, x_0, 0+)
$$

and calculate

$$
e'(r) = -\frac{n+2\kappa-2}{r} (e(r) + W(\mathbf{u}, x_0, 0+)) + \kappa r^{-n-2\kappa} \int_{\partial B_r(x_0)} |\mathbf{u}|^2 d\mathcal{H}^{n-1}
$$

+ $r^{-n-2\kappa+2} \int_{\partial B_r(x_0)} (|\nabla \mathbf{u}|^2 + 2F(\mathbf{u})) d\mathcal{H}^{n-1}$
- $2\kappa r^{-n-2\kappa+1} \int_{\partial B_r(x_0)} (\nabla \mathbf{u} \cdot \nu) \cdot \mathbf{u} d\mathcal{H}^{n-1} - \kappa (n-1) r^{-n-2\kappa} \int_{\partial B_r(x_0)} |\mathbf{u}|^2 d\mathcal{H}^{n-1}$
= $-\frac{n+2\kappa-2}{r} (e(r) + W(\mathbf{u}, x_0, 0+)) - \frac{\kappa}{r} (n-2) \int_{\partial B_1} |\mathbf{u}_r|^2 d\mathcal{H}^{n-1}$
+ $\frac{1}{r} \int_{\partial B_1} (|\nabla \mathbf{u}_r|^2 + 2F(\mathbf{u}_r)) d\mathcal{H}^{n-1} - \frac{2\kappa}{r} \int_{\partial B_1} (\nabla \mathbf{u}_r \cdot \nu) \cdot \mathbf{u}_r d\mathcal{H}^{n-1}$
 $\geq -\frac{n+2\kappa-2}{r} (e(r) + W(\mathbf{u}, x_0, 0+))$
+ $\frac{1}{r} \int_{\partial B_1} |\nabla_\theta \mathbf{u}_r|^2 - (\kappa (n-2) + \kappa^2) |\mathbf{u}_r|^2 + 2F(\mathbf{u}_r) d\mathcal{H}^{n-1}$
= $-\frac{n+2\kappa-2}{r} (e(r) + W(\mathbf{u}, x_0, 0+))$
+ $\frac{1}{r} \int_{\partial B_1} |\nabla_\theta \mathbf{c}_r|^2 - (\kappa (n-2) + \kappa^2) |\mathbf{c}_r|^2 + 2F(\mathbf{c}_r) d\mathcal{H}^{n-1}$
= $-\frac{n+2\kappa-2}{r} (e(r) + W(\mathbf{u}, x_0, 0+$

At this point, we employ the minimality of u as well as the assumption that the epiperimetric inequality $M(\mathbf{v}) \leq (1 - \varepsilon)M(\mathbf{c}_r) + \varepsilon W(\mathbf{u}, x_0, 0+)$ holds for some $\mathbf{v} \in W^{1,2}(B_1;\mathbb{R}^m)$ with \mathbf{c}_r -boundary values and we obtain for $r \in (0,r_0)$ the estimate

$$
e'(r) \geq \frac{n+2\kappa-2}{r} \left(\frac{1}{1-\varepsilon} (M(\mathbf{u}_r) - W(\mathbf{u}, x_0, 0+)) - e(r) \right)
$$

=
$$
\frac{n+2\kappa-2}{r} \left(\frac{1}{1-\varepsilon} - 1 \right) e(r) = \frac{n+2\kappa-2}{r} \frac{\varepsilon}{1-\varepsilon} e(r).
$$

By the monotonicity formula Proposition [2.3,](#page-4-0) $e(r) \geq 0$, and we conclude in the non-trivial case $e > 0$ that in (r_1, r_0)

$$
e(r)\leq e(r_0)\big(\frac{r}{r_0}\big)^{\Lambda}\quad\text{for }r\in(r_1,r_0),
$$

which proves the first statement.

Now using once more the monotonicity formula, Proposition [2.3,](#page-4-0) we get for $0 < \rho < \sigma \leq r_0$ an estimate of the form

$$
\int_{\partial B_1} |\mathbf{u}_{\sigma}(x) - \mathbf{u}_{\rho}(x)| d\mathcal{H}^{n-1} \leq \int_{\partial B_1} \int_{\rho}^{\sigma} |\partial_r \mathbf{u}_r| dr d\mathcal{H}^{n-1}
$$
\n
$$
= \int_{\rho}^{\sigma} r^{-1} \int_{\partial B_1} |\nabla \mathbf{u}_r \cdot x - \kappa \mathbf{u}_r| d\mathcal{H}^{n-1} dr
$$
\n
$$
\leq \int_{\rho}^{\sigma} r^{-1/2} \sqrt{\frac{n\omega_n}{2}} \Big(\int_{\partial B_1} \frac{2}{r} |\nabla \mathbf{u}_r \cdot x - \kappa \mathbf{u}_r|^2 d\mathcal{H}^{n-1} \Big)^{1/2} dr
$$
\n
$$
= \sqrt{\frac{n\omega_n}{2}} \int_{\rho}^{\sigma} r^{-1/2} \sqrt{e'(r)} dr
$$
\n
$$
\leq \sqrt{\frac{n\omega_n}{2}} (\log(\sigma) - \log(\rho))^{1/2} (e(\sigma) - e(\rho))^{1/2}.
$$

Considering now $0 < 2\rho < 2r \le r_0$ and intervals $[2^{-k-1}, 2^{-k}) \ni \rho$ and $[2^{-\ell-1}, 2^{-\ell}) \ni$ r the already proved part of the theorem yields that

$$
\int_{\partial B_1} |\mathbf{u}_r(x) - \mathbf{u}_\rho(x)| d\mathcal{H}^{n-1} \le C_1(n) \sum_{i=\ell}^k (\log(2^{-i}) - \log(2^{-i-1}))^{1/2} (e(2^{-i}) - e(2^{-i-1}))^{1/2}
$$
\n
$$
\le C_2(n) \sum_{i=\ell}^k (e(2^{-i}) - e(2^{-i-1}))^{1/2}
$$
\n
$$
\le C_2(n) (e(r_0))^{1/2} \sum_{i=\ell}^k (r_0 2^i)^{-\Lambda/2}
$$
\n
$$
\le C_2(n) (e(r_0))^{1/2} r_0^{-\Lambda/2} \sum_{i=\ell}^\infty 2^{-i\Lambda/2}
$$
\n
$$
\le C_3(n, \kappa, \varepsilon) (e(r_0))^{1/2} (r_0 2^\ell)^{-\Lambda/2}
$$
\n
$$
\le C_3(n, \kappa, \varepsilon) (e(r_0))^{1/2} (\frac{2r}{r_0})^{\Lambda/2}.
$$

4.3. Uniqueness of blow-up. In this subsection, we will show the uniqueness of blow-up and estimate the rate of convergence of the scaled solution to its blow-up in Theorem [4.7.](#page-16-1)

First we need to prove some preliminaries.

Lemma 4.3. The half-plane solutions of the system [\(3\)](#page-1-3) are isolated (in the topology of $W^{1,2}(B_1(0); \mathbb{R}^m)$) within the class of homogeneous solutions of degree κ .

Proof. We suppose towards a contradiction that this does not hold. Then there exists a sequence of homogeneous solutions of degree κ , say \mathbf{u}_n , such that

$$
0 < \inf_{\mathbf{h}\in\mathbb{H}} \|\mathbf{u}_n - \mathbf{h}\|_{W^{1,2}(B_1(0);\mathbb{R}^m)} = \|\mathbf{u}_n - \hat{\mathbf{h}}\|_{W^{1,2}(B_1(0);\mathbb{R}^m)} =: \delta_n \to 0, \quad \text{as } n \to \infty,
$$

where $\hat{\mathbf{h}} = \alpha(x_n^+)^{\kappa} \mathbf{e}_1$. When passing to a subsequence, $(\mathbf{u}_n - \hat{\mathbf{h}})/\delta_n =: \mathbf{w}_n \rightharpoonup \mathbf{w}$ weakly in $W^{1,2}(B_1(0);\mathbb{R}^m)$, the limit w is still a homogeneous function of degree *κ*. Furthermore, for $\phi \in C_0^{\infty}(B_1; \mathbb{R}^m)$ we have

$$
-\int_{B_1} \nabla \mathbf{w}_n \cdot \nabla \phi \, dx = \frac{1}{\delta_n} \int_{B_1} (f(\mathbf{u}_n) - f(\hat{\mathbf{h}})) \cdot \phi \, dx
$$

$$
= \frac{1}{\delta_n} \int_{B_1} \int_0^1 \frac{d}{dt} f(\hat{\mathbf{h}} + t(\mathbf{u}_n - \hat{\mathbf{h}})) \cdot \phi \, dt dx
$$

$$
= \int_{B_1} \int_0^1 f_\mathbf{u}(\hat{\mathbf{h}} + t\delta_n \mathbf{w}_n)(\mathbf{w}_n) \cdot \phi \, dt dx.
$$

If supp $\phi \subset B_1^-$, let $n \to \infty$ we conclude that

$$
\int_{B_1^-} f_{\mathbf{u}}(\mathbf{w})(\mathbf{w}) \cdot \phi \, dx = \lim_{n \to \infty} \int_{B_1^-} f_{\mathbf{u}}(\mathbf{w}_n)(\mathbf{w}_n) \cdot \phi \, dx
$$

$$
= - \lim_{n \to \infty} \int_{B_1^-} q \delta_n^{1-q} \nabla \mathbf{w}_n \cdot \nabla \phi \, dx = 0.
$$

Then $\mathbf{w} \equiv 0$ in $B_1(0) \cap \{x_n < 0\}$. Now for every supp $\phi \subset B_1^+$,

$$
-\int_{B_1^+} \nabla \mathbf{w} \cdot \nabla \phi \, dx = \int_{B_1^+} f_{\mathbf{u}}(\hat{\mathbf{h}})(\mathbf{w}) \cdot \phi \, dx.
$$

Thus $\Delta \mathbf{w} = f_{\mathbf{u}}(\hat{\mathbf{h}})(\mathbf{w})$ in $B_1(0) \cap \{x_n > 0\}.$

On the other hand, we know that \hat{h} is the best approximation to u_n among all half-plane solutions. But then it follows exactly as in Step 3 of the proof of the epiperimetric inequality Theorem [3.1](#page-6-1) that $w \equiv 0$. In order to obtain a contradiction to the assumption $\delta_n > 0$ by which $\|\mathbf{w}_n\|_{W^{1,2}(B_1(0);\mathbb{R}^m)} = 1$, it is therefore sufficient to show the strong convergence of ∇ **w**_n to ∇ **w** in $\mathcal{L}^2(B_1(0); \mathbb{R}^{mn})$ as the subsequence $n \to \infty$. But by compact imbedding on the boundary

$$
\int_{B_1} |\nabla \mathbf{w}_n|^2 = \int_{\partial B_1} \mathbf{w}_n \cdot \nabla \mathbf{w}_n \nu d\mathcal{H}^{n-1} - \int_{B_1} \mathbf{w}_n \cdot \Delta \mathbf{w}_n
$$

\n
$$
= \kappa \int_{\partial B_1} |\mathbf{w}_n|^2 d\mathcal{H}^{n-1} - \frac{1}{\delta_n^2} \int_{B_1} (\mathbf{u}_n - \hat{\mathbf{h}}) \cdot (f(\mathbf{u}_n) - f(\hat{\mathbf{h}})) dx
$$

\n
$$
\leq \kappa \int_{\partial B_1} |\mathbf{w}_n|^2 d\mathcal{H}^{n-1} \to 0,
$$

as the subsequence $n \to \infty$.

Remark 4.4. Theorem [4.2](#page-11-3) proves the uniqueness of blow-ups provided \mathbf{u}_r remains in a δ -neighborhood of \mathbb{H} , where δ is the constant introduced in the epiperimetric inequality. Lemma [4.3,](#page-13-1) however, provides this condition.

Proposition 4.5. Let $u \neq 0$ be a homogeneous solution of degree κ satisfying ${||u| = 0}^{\circ} \neq \emptyset$. Then $M(u) \ge \alpha_n/2$, and equality implies that **u** is a half-plane solution; here $\alpha_n = 2M(\mathbf{h})$ for every $\mathbf{h} \in \mathbb{H}$.

Proof. The proof is by induction on n , the dimension of the domain. In one space dimension the statement is an immediate consequence of the homogeneity. We assume that it holds for every solution in dimension $\leq n-1$ and that it is violated by a homogeneous solution **u** of degree κ in dimension n, that $\{|{\bf u}| = 0\}$ contains the ball B and that $\mathbf{e}_n \in \partial B \cap \partial {\{\mathbf{u}\}\n}$. The homogeneity of **u** implies that

$$
W(\mathbf{u},\mathbf{e}_n,0+) = \lim_{r \to 0^+} W(\mathbf{u},\mathbf{e}_n,r) = \lim_{r \to 0^+} W(\mathbf{u},\frac{\mathbf{e}_n}{m},\frac{r}{m}) = W(\mathbf{u},\frac{\mathbf{e}_n}{m},0+),
$$

and by the upper semicontinuity of the function $x \mapsto W(\mathbf{u}, x, 0+)$,

$$
W(\mathbf{u},\mathbf{e}_n,0+) = \limsup_{m\to\infty} W(\mathbf{u},\frac{\mathbf{e}_n}{m},0+) \leq W(\mathbf{u},0,0+) \leq W(\mathbf{u},0,1) = M(\mathbf{u}) < \frac{\alpha_n}{2}.
$$

Thus every blow-up limit \mathbf{u}_0 of \mathbf{u} at the point \mathbf{e}_n satisfies the inequality $M(\mathbf{u}_0)$ < $\alpha_n/2$. (Note that by the nondegeneracy property $\mathbf{u}_0 \neq 0$.) Now the homogeneity of **u** tells us that \mathbf{u}_0 must be constant in the direction of the vector \mathbf{e}_n and that again $\{|{\bf u}_0|=0\}^\circ \neq \emptyset$, so $\tilde{\bf u}:={\bf u}_0|_{\mathbb{R}^{n-1}}$ is a homogeneous solution of degree κ satisfying $\{|\tilde{\mathbf{u}}|=0\}^{\circ}\neq\varnothing$, and

$$
\frac{\alpha_n}{2} > \int_{B_1} |\nabla \mathbf{u}_0|^2 + 2F(\mathbf{u}_0) dx - \kappa \int_{\partial B_1} |\mathbf{u}_0|^2 d\mathcal{H}^{n-1} = \frac{1-q}{1+q} \int_{B_1} |\mathbf{u}_0|^{1+q} dx
$$

\n
$$
= \frac{2(1-q)}{1+q} \int_{\{|x'|<1\}} \int_0^{\sqrt{1-|x'|^2}} |\tilde{\mathbf{u}}(x')|^{1+q} dx_n dx'
$$

\n
$$
= \frac{2(1-q)}{1+q} \int_0^1 \sqrt{1-r^2} r^{n-2+\kappa(1+q)} dr \int_{\partial B'_1} |\tilde{\mathbf{u}}(x')|^{1+q} d\mathcal{H}^{n-2}
$$

\n
$$
= 2M(\tilde{\mathbf{u}}) \int_0^1 \sqrt{1-r^2} r^{n-2+\kappa(1+q)} dr
$$

\n
$$
\geq \alpha_{n-1} \int_0^1 \sqrt{1-r^2} r^{n-2+\kappa(1+q)} dr = \frac{\alpha_n}{2},
$$

which contradicts the induction hypothesis. (Notice that the last equality is obtained by the same calculation for $h(x) = \alpha(x_1^+)^{\kappa} e \in \mathbb{H}$ instead of u_0 .

Finally, we assume inductively that the second part of the statement holds for every dimension $\leq n-1$ and consider the case of a homogeneous solution **u** of degree κ in dimension n satisfying $M(\mathbf{u}) = \alpha_n/2$, $B \subset \{|\mathbf{u}| = 0\}$ and $\mathbf{e}_n \in \partial B \cap \partial {\{|\mathbf{u}| > 0\}}$. As in the first part of the proof we obtain that every blow-up limit \mathbf{u}_0 of \mathbf{u} at the point \mathbf{e}_n satisfies the inequlity $M(\mathbf{u}_0) \leq \alpha_n/2$, that \mathbf{u}_0 is constant in the direction of \mathbf{e}_n and that $\{|\mathbf{u}_0|=0\}^\circ \neq \varnothing$. Defining again $\tilde{\mathbf{u}} := \mathbf{u}_0|_{\mathbb{R}^{n-1}}$, which is a homogeneous solution of degree κ satisfying $\{|\tilde{\mathbf{u}}| = 0\}^{\circ} \neq \emptyset$, the calculation in the first part of the proof yields that $M(\tilde{u}) \leq \alpha_{n-1}/2$. Thus \tilde{u} must be a half-plane solution by the induction hypothesis, and so must \mathbf{u}_0 . Therefore, for $0 < r_m \to 0$, every blow-up limit of **u** at the point $r_m \mathbf{e}_n$ must be a half-plane solution. Assuming that $\mathbf{u} \notin \mathbb{H}$, we find by a continuity argument for an arbitrary $\theta \in (0,1)$ a sequence $\rho_m \to 0$ such that

$$
dist(\rho_m^{-\kappa} \mathbf{u}(r_m \mathbf{e}_n + \rho_m \cdot), \mathbb{H}) = \theta dist(\mathbf{u}, \mathbb{H}) > 0,
$$

where the distance is measured in the $W^{1,2}(B_1(0); \mathbb{R}^m)$ -norm. On the other hand, it follows that $\mathbf{u}(r_m \mathbf{e}_n + \rho_m \cdot)/\rho_m^{\kappa}$ converges in $W^{1,2}(B_1(0); \mathbb{R}^m)$ to a homogeneous solution \mathbf{u}^* of degree κ along a subsequence $m \to \infty$. Although the boundedness of $\mathbf{u}(r_m \mathbf{e}_n + \rho_m \cdot)/\rho_m^{\kappa}$ implies the weak convergence in $W^{1,2}(B_1(0); \mathbb{R}^m)$, the compact embedding on the boundary proves the strong convergence as at the end of the proof of Lemma [4.3.](#page-13-1) The conclusion is that $dist(\mathbf{u}^*, \mathbb{H}) = \theta dist(\mathbf{u}, \mathbb{H}) > 0$ which for small θ contradicts the isolation property Lemma [4.3.](#page-13-1)

Remark 4.6. Proposition [4.5](#page-14-1) and Lemma [4.3](#page-13-1) show that the infimum energy of all κ -homogeneous solutions outside of $\mathbb H$ is strictly greater than $\alpha_n/2$. From this fact we infer that the set of regular free boundary points $\mathcal{R}_{\mathbf{u}}$ is open relative to $\Gamma(\mathbf{u})$.

Theorem 4.7. Let C_h be a compact subset of \mathcal{R}_u . Assume that $\mathbf{u}_0(x) = \alpha \max(x \cdot$ $\nu(x_0), 0^k e(x_0)$ is the blow-up limit of **u** at x_0 which is a half-plane solution, for some $\nu(x_0) \in \partial B_1(0) \subset \mathbb{R}^n$ and $e(x_0) \in \partial B_1(0) \subset \mathbb{R}^m$. Then there exists $r_0 > 0$ and positive constant C, such that

$$
\int_{\partial B_1} \left| \frac{\mathbf{u}(x_0 + rx)}{r^{\kappa}} - \alpha \max(x \cdot \nu(x_0), 0)^{\kappa} \mathbf{e}(x_0) \right| d\mathcal{H}^{n-1} \le Cr^{\Lambda/2},
$$

for every $x_0 \in C_h$ and every $r \leq r_0$. Here, Λ is the exponent defined in Theorem [4.2.](#page-11-3)

Proof. In view of Theorem [3.1](#page-6-1) and Theorem [4.2,](#page-11-3) it is sufficient to show that

$$
dist(\frac{\mathbf{u}(x_0+r\,\cdot)}{r^{\kappa}}, \mathbb{H}) \le \delta,
$$

for every $x_0 \in C_h$ and $r \leq r_0$, where the distance is measured in the $W^{1,2}(B_1(0); \mathbb{R}^m)$ norm. Assume now towards a contradiction that $dist(\frac{u(x_i+\rho_i \cdot)}{\rho_i^{\kappa}}, \mathbb{H}) \ge \delta > 0$ for some $x_i \in C_h$ and $\rho_i \to 0$. By a continuity argument, for each $\theta \in (0,1)$ there is a sequence $\tilde{\rho}_i < \rho_i$ such that $dist(\frac{\mathbf{u}(x_i + \tilde{\rho}_i \cdot)}{\tilde{\rho}_i^{\kappa}}, \mathbb{H}) = \theta \delta$. Then $\mathbf{u}_i := \frac{\mathbf{u}(x_i + \tilde{\rho}_i \cdot)}{\tilde{\rho}_i^{\kappa}}$ is bounded in $W^{1,2}(B_1(0); \mathbb{R}^m)$, and passing to a limit with respect to a subsequence we obtain a solution \mathbf{u}_h satisfying $dist(\mathbf{u}_h, \mathbb{H}) = \theta \delta$. Moreover,

$$
W(\mathbf{u}_h, 0, r) = \lim_{i \to \infty} W(\mathbf{u}_i, 0, r) = \lim_{i \to \infty} W(\mathbf{u}, x_i, r\tilde{\rho}_i) = \alpha_n/2.
$$

Thus \mathbf{u}_h is a κ -homogeneous solution by Proposition [2.3](#page-4-0) and so for small θ contra-dicts the isolation property in Lemma [4.3.](#page-13-1)

4.4. **Proof of Theorem [1.4.](#page-2-1)** Let us consider $x_0 \in \mathcal{R}_u$. By Theorem [4.7](#page-16-1) there exists $\delta_0 > 0$ such that $B_{2\delta_0}(x_0) \subset B_1$, $B_{2\delta_0}(x_0) \cap \partial\{|{\bf u}| > 0\} = B_{2\delta_0}(x_0) \cap \mathcal{R}_{\bf u}$ and

(19)
$$
\int_{\partial B_1} \left| \frac{\mathbf{u}(x_1+rx)}{r^{\kappa}} - \alpha \max(x \cdot \nu(x_1), 0)^{\kappa} \mathbf{e}(x_1) \right| d\mathcal{H}^{n-1} \leq Cr^{\Lambda/2},
$$

for every $x_1 \in \partial\{|u| > 0\} \cap B_{\delta_0}(x_0)$ and for every $r \leq \min(\delta_0, r_0)$. We now observe that $x_1 \mapsto \nu(x_1)$ and $x_1 \mapsto e(x_1)$ are Hölder continuous with exponent β on $\partial\{|{\bf u}| > 0\} \cap B_{\delta_1}(x_0)$ for some $\delta_1 \in (0, \delta_0)$:

$$
\alpha \int_{\partial B_1} \left| \max(x \cdot \nu(x_1), 0)^{\kappa} \mathbf{e}(x_1) - \max(x \cdot \nu(x_2), 0)^{\kappa} \mathbf{e}(x_2) \right| d\mathcal{H}^{n-1}
$$

\n
$$
\leq 2Cr^{\Lambda/2} + \int_{\partial B_1} \int_0^1 \left| \frac{\nabla \mathbf{u}(x_1 + rx + t(x_2 - x_1))}{r^{\kappa}} \right| |x_1 - x_2| dt d\mathcal{H}^{n-1}
$$

\n
$$
\leq 2Cr^{\Lambda/2} + C_1 \frac{|x_1 - x_2|}{r^{\kappa}} \leq (2C + C_1)|x_1 - x_2|^{\gamma \Lambda/2},
$$

if we choose $\gamma := (\kappa + \frac{\Lambda}{2})^{-1}$ and $r := |x_1 - x_2|^{\gamma} \le \min(\delta_0, r_0)$, and the left hand side

(20)
$$
\alpha \int_{\partial B_1} \left| \max(x \cdot \nu(x_1), 0)^{\kappa} \mathbf{e}(x_1) - \max(x \cdot \nu(x_2), 0)^{\kappa} \mathbf{e}(x_2) \right| d\mathcal{H}^{n-1} \ge c(n) \left(|\nu(x_1) - \nu(x_2)| + |\mathbf{e}(x_1) - \mathbf{e}(x_2)| \right)
$$

as can be shown by an indirect argument. Suppose towards a contradiction that $c_j := |\nu_j^1 - \nu_j^2| + |\mathbf{e}_j^1 - \mathbf{e}_j^2| \to 0, \, (\nu_j^1 - \nu_j^2)/c_j \to \eta, \, (\mathbf{e}_j^1 - \mathbf{e}_j^2)/c_j \to \xi, \, \nu_j^1 \to \bar{\nu}, \, \mathbf{e}_j^2 \to \bar{\mathbf{e}}$

and
\n
$$
0 \leftarrow \frac{1}{c_j} \int_{\partial B_1} \left| \max(x \cdot \nu_j^1, 0)^{\kappa} \mathbf{e}_j^1 - \max(x \cdot \nu_j^2, 0)^{\kappa} \mathbf{e}_j^2 \right| d\mathcal{H}^{n-1}
$$
\n
$$
\geq \frac{1}{c_j} \int_{\partial B_1 \cap \{x \cdot \nu_j^1 > 0\} \cap \{x \cdot \nu_j^2 > 0\}} \left| (x \cdot \nu_j^1)^{\kappa} (\mathbf{e}_j^1 - \mathbf{e}_j^2) - \left((x \cdot \nu_j^1)^{\kappa} - (x \cdot \nu_j^2)^{\kappa} \right) \mathbf{e}_j^2 \right| d\mathcal{H}^{n-1}
$$
\n
$$
\to \int_{\partial B_1 \cap \{x \cdot \bar{\nu} > 0\}} \left| (x \cdot \bar{\nu})^{\kappa} \xi - \kappa (x \cdot \bar{\nu})^{\kappa - 1} (x \cdot \eta) \bar{\mathbf{e}} \right| d\mathcal{H}^{n-1}.
$$

Then $(x \cdot \bar{\nu})\xi = \kappa(x \cdot \eta)$ **e** for all $x \in \partial B_1 \cap \{x \cdot \bar{\nu} > 0\}$. Putting $x = \bar{\nu}$ and noticing that $0 = (|\nu_j^1|^2 - |\nu_j^2|^2)/c_j \to 2\eta \cdot \bar{\nu}$, it follows that $\xi = 0$. So, $x \cdot \eta = 0$ for all x implying that $\eta = 0$. But $|\xi| + |\eta| = 1$. This contradiction proves [\(20\)](#page-16-2).

Next, [\(19\)](#page-16-3) as well as the regularity and nondegenracy of **u** imply that for $\epsilon > 0$ there exists $\delta_2 \in (0, \delta_1)$ such that for $x_1 \in \partial\{|{\bf u}| > 0\} \cap B_{\delta_1}(x_0)$ and $y \in B_{\delta_2}(x_1)$,

(21)
$$
\mathbf{u}(y) = 0 \quad \text{if} \quad (y - x_1) \cdot \nu(x_1) < -\epsilon |y - x_1|, |\mathbf{u}(y)| > 0 \quad \text{if} \quad (y - x_1) \cdot \nu(x_1) > \epsilon |y - x_1|.
$$

Assuming that [\(21\)](#page-17-1) does not hold, we obtain a sequence $\partial\{|{\bf u}| > 0\} \cap B_{\delta_1}(x_0) \ni$ $x_m \to \bar{x}$ and a sequence $y_m - x_m \to 0$ as $m \to \infty$ such that

(22) either
$$
\mathbf{u}(y_m) = 0
$$
 and $(y_m - x_m) \cdot \nu(x_m) < -\epsilon |y_m - x_m|$,
or $|\mathbf{u}(y_m)| > 0$ and $(y_m - x_m) \cdot \nu(x_m) > \epsilon |y_m - x_m|$.

On the other hand we know from [\(19\)](#page-16-3) as well as from the regularity and nondegeneracy of the solution **u**, that the sequence $\mathbf{u}_j(x) := \frac{\mathbf{u}(x_j + |y_j - x_j|x)}{|y_j - x_j|^{\kappa}}$ converges in $C^{1,\alpha}_{loc}(\mathbb{R}^n;\mathbb{R}^m)$ to $\alpha \max(x \cdot \nu(\bar{x}),0)^\kappa \mathbf{e}(\bar{x})$ as $j \to \infty$ and that $\mathbf{u}_j = 0$ on each compact subset of $\{x \cdot \nu(\bar{x}) < 0\}$ provided that $j \geq j(C)$. This, however, contradicts [\(22\)](#page-17-2) for large j.

Last, we use [\(21\)](#page-17-1) in order to show that $\partial\{|{\bf u}| > 0\}$ is for some $\delta_3 \in (0, \delta_2)$ in $B_{\delta_3}(x_0)$ the graph of a differentiable function. Applying two rotations we may assume that $\nu(x_0) = \mathbf{e}_n$ and $\mathbf{e}(x_0) = \mathbf{e}_1$. Choosing now δ_2 with respect to $\epsilon = \frac{1}{2}$ and defining functions $g^+, g^- : B'_{\frac{\delta_2}{2}}(0) \to [-\infty, \infty],$

$$
g^{+}(x') := \sup\{x_n : x_0 + (x', x_n) \in \partial\{|u| > 0\}\},\
$$
and

$$
g^{-}(x') := \inf\{x_n : x_0 + (x', x_n) \in \partial\{|u| > 0\}\},\
$$

we infer from [\(21\)](#page-17-1) as well as from the continuity of $\nu(x)$ immediately that $\{x_n:$ $x_0 + (x', x_n) \in \partial\{|u| > 0\}$ is non-empty and that for sufficiently small δ_3 the functions g^+ and g^- are Lipschitz continuous and satisfy $g^+ = g^-$ on $\overline{B'_{\delta_3}(0)}$. Applying [\(21\)](#page-17-1) once more with respect to arbitrary ϵ we see that g^+ is Fréchetdifferentiable in $\overline{B'_{\delta_3}(0)}$, which finishes our proof in view of the already derived Hölder continuity of the normal vector $\nu(x)$.

5. A SYSTEM WITH HÖLDER COEFFICIENTS

In this section we are going to show that our results extend to the setting

(23)
$$
\Delta \mathbf{u} = f(x, \mathbf{u}) := \lambda_+(x) |\mathbf{u}^+|^{q-1} \mathbf{u}^+ - \lambda_-(x) |\mathbf{u}^-|^{q-1} \mathbf{u}^-, \quad \text{in } B_1(0),
$$

where $\mathbf{u}^{\pm} = (u_1^{\pm}, \cdots, u_m^{\pm})$ and $u_i^{\pm} := \max(\pm u_i, 0)$. Also, we assume that coefficients λ_{\pm} satisfy

$$
0 < \lambda_0 \leq \lambda_{\pm} \in C^{0,\beta}(B_1).
$$

In fact, u is a minimizer of

$$
J(\mathbf{u}) := \int_{B_1} |\nabla \mathbf{u}|^2 + 2F(x, \mathbf{u}) dx,
$$

where $F(x, \mathbf{u}) := \frac{1}{1+q} (\lambda_+(x)|\mathbf{u}^+|^{q+1} + \lambda_-(x)|\mathbf{u}^-|^{q+1}).$

First of all, we see that first part of Theorem [1.2](#page-2-4) (case $\kappa \notin \mathbb{N}$) is valid. Even in the case $\kappa \in \mathbb{N}$, we may use its result to obtain the estimate

$$
\sup_{B_r(z)} |\mathbf{u}| \le C_{\varepsilon} r^{\kappa-\varepsilon},
$$

for every $\varepsilon > 0$. (Remark [2.4\)](#page-6-5) It implies the monotonicity, Proposition [2.3,](#page-4-0) for the energy

$$
W_s(\mathbf{u}, x_0, r) = \frac{1}{r^{n+2\kappa - 2}} \int_{B_r(x_0)} \left(|\nabla \mathbf{u}|^2 + 2F_s(x, \mathbf{u}) \right) dx - \frac{\kappa}{r^{n+2\kappa - 1}} \int_{\partial B_r(x_0)} |\mathbf{u}|^2 d\mathcal{H}^{n-1}
$$

where $F_s(x, \mathbf{u}) := F(x_0 + s(x - x_0), \mathbf{u})$. We have the relation $W_s(r, x_0, \mathbf{u}_r) =$ $W_1(rs, x_0, u)$ and the following (almost) monotonicity. (see [\[5\]](#page-21-11) for similar setting but in the scalar case.)

Proposition 5.1. Let **u** be a solution of [\(23\)](#page-17-3) in $B_{r_0}(x_0)$ and $x_0 \in \Gamma^{\kappa}(\mathbf{u})$. There exist constants $C > 0$ and $\mu > 0$ such that $W_1(\mathbf{u}, x_0, r) + c r^{\mu}$ is increasing for $r > 0$.

By the monotonicity we can repeat the second part of the proof of Theorem [1.2.](#page-2-4) Moreover, we have still non-degeneracy property as we have shown in Proposition [4.1.](#page-11-4) When coefficients λ_{\pm} are constant, we can repeat the proof of Theorem [3.1](#page-6-1) to show the Epiperimetric inequality for the energy function

$$
M_{x_0}(\mathbf{v}) := \int_{B_1} \left(|\nabla \mathbf{v}|^2 + 2F(x_0, \mathbf{v}) \right) dx - \kappa \int_{\partial B_1} |\mathbf{v}|^2 d\mathcal{H}^{n-1}.
$$

We now proceed with the proof of the regularity of free boundary at regular points.

Theorem 5.2 (Energy decay). Let $x_0 \in B_1 \cap \partial\{|u| > 0\}$, and suppose that the epiperimetric inequality holds with $\varepsilon \in (0,1)$ for each

$$
\mathbf{c}_r(x) := |x|^{\kappa} \mathbf{u}_r(\frac{x}{|x|}) = \frac{|x|^{\kappa}}{r^{\kappa}} \mathbf{u}(x_0 + \frac{r}{|x|}x)
$$

and for all $r \le r_0 < 1$. Finally let \mathbf{u}_0 denote an arbitrary blow-up limit of \mathbf{u} at x_0 and $\Lambda = \min\{(n + 2\kappa - 2)\varepsilon/(1 - \varepsilon), \beta\}$. Then there exists a constant C such that

$$
|W_1(\mathbf{u}, x_0, r) - W_1(\mathbf{u}, x_0, 0+)| \le Cr^{\Lambda} |\log r|,
$$

and

$$
\int_{\partial B_1} |\mathbf{u}_r(x) - \mathbf{u}_0(x)| d\mathcal{H}^{n-1} \le Cr^{\Lambda/2} |\log r|,
$$

for all $r \in (0, r_0)$.

Proof. We can repeat the calculation in the proof of Theorem [4.2](#page-11-3) to show that

$$
e'(r) \geq -\frac{n+2\kappa-2}{r} \big(e(r) + W_1(\mathbf{u}, x_0, 0+) - W_r(\mathbf{c}_r, x_0, 1)\big).
$$

Now we apply the epiperimetric inequality to \mathbf{c}_r and find a function $\mathbf{v} \in W^{1,2}(B_1;\mathbb{R}^m)$ such that

$$
M_{x_0}(\mathbf{v}) \le (1 - \varepsilon) M_{x_0}(\mathbf{c}_r) + \varepsilon M_{x_0}(\mathbf{h}).
$$

,

Moreover, we may assume that

$$
(24) \t\t\t M_{x_0}(\mathbf{v}) \leq M_{x_0}(\mathbf{u}_r),
$$

otherwise we substitute **v** by \mathbf{u}_r . In order to replace **v** by \mathbf{u}_r generally, we find the following estimate by the Hölder regularity assumption on $F(x, v)$ with respect to the variable x:

(25)
$$
|M_{x_0}(\mathbf{v}) - W_r(\mathbf{v}, x_0, 1)| \leq C_1 r^{\beta} ||\mathbf{v}||_{\mathcal{L}^{1+q}(B_1)}^{1+q},
$$

for some constant C depending only on the coefficients of the problem. Freezing the coefficients and estimate [\(25\)](#page-19-0) yields that

$$
M_{x_0}(\mathbf{v}) \ge W_r(\mathbf{v}, x_0, 1) - C_1 r^{\beta} ||\mathbf{v}||_{\mathcal{L}^{1+q}(B_1)}^{1+q}
$$

\n
$$
\ge W_r(\mathbf{v}, x_0, 1) - \frac{1}{2} C_1 r^{\beta} (M_{x_0}(\mathbf{v}) + \kappa ||\mathbf{v}||_{\mathcal{L}^2(\partial B_1)}^2)
$$

\n
$$
\ge W_r(\mathbf{v}, x_0, 1) - \frac{1}{2} C_1 r^{\beta} (M_{x_0}(\mathbf{u}_r) + \kappa ||\mathbf{u}_r||_{\mathcal{L}^2(\partial B_1)}^2),
$$

where we have used [\(24\)](#page-19-1) in the last line. Now by the minimality of \mathbf{u}_r with the respect of its boundary conditions, we have that

$$
e'(r) \geq -\frac{n+2\kappa-2}{r} \left(e(r) + M_{x_0}(\mathbf{h}) - M_{x_0}(\mathbf{c}_r) + C_1 r^{\beta} ||\mathbf{c}_r||_{\mathcal{L}^{1+q}(B_1)}^{1+q} \right)
$$

\n
$$
\geq -\frac{n+2\kappa-2}{r} \left(e(r) + \frac{M_{x_0}(\mathbf{h}) - M_{x_0}(\mathbf{v})}{1-\varepsilon} + C_1 r^{\beta} ||\mathbf{c}_r||_{\mathcal{L}^{1+q}(B_1)}^{1+q} \right)
$$

\n
$$
\geq -\frac{n+2\kappa-2}{r} \left(e(r) + \frac{1}{1-\varepsilon} \left[M_{x_0}(\mathbf{h}) - W_r(\mathbf{v}, x_0, 1) + \frac{1}{2} C_1 r^{\beta} \left(M_{x_0}(\mathbf{u}_r) + \kappa ||\mathbf{u}_r||_{\mathcal{L}^2(\partial B_1)}^2 \right) \right] + C_1 r^{\beta} ||\mathbf{c}_r||_{\mathcal{L}^{1+q}(B_1)}^{1+q} \right)
$$

\n
$$
\geq -\frac{n+2\kappa-2}{r} \left[e(r) + \frac{1}{1-\varepsilon} \left(M_{x_0}(\mathbf{h}) - W_r(\mathbf{u}_r, x_0, 1) \right) + C_2 r^{\beta} \right]
$$

\n
$$
= \frac{n+2\kappa-2}{r} \frac{\varepsilon}{1-\varepsilon} e(r) - C_3 r^{\beta-1} \geq \Lambda e(r) - C_3 r^{\Lambda-1}.
$$

Therefore

$$
r^{-\Lambda}e(r))' \geq -C_3r^{-1},
$$

 $\left($

and after integrating in (r, r_0) , we get

$$
e(r) \le Cr^{\Lambda} |\log r|.
$$

The second part of theorem can be obtained similarly to the proof of Theorem $4.2.$

Remark 5.3. As apparent from the proof, we need the term $\log r$ in the estimate in Theorem [5.2](#page-18-0) only when $\beta = (n + 2\kappa - 2)\varepsilon/(1 - \varepsilon)$. Otherwise, the estimations are valid without this term.

Other results in Section 4 remain true when we replace the equation with [\(23\)](#page-17-3), especially Theorem [1.4](#page-2-1) on the regularity of the free boundary.

Remark 5.4. In this section we can treat case $q = 0$ (or $\kappa = 2$). Although, the proof of epiperimetric inequality in Section [3](#page-6-0) needs the condition $0 < q < 1$, we already know the similar result in [\[2\]](#page-21-1). Then our proof in Theorem [5.2](#page-18-0) covers the case $q=0$.

6. Appendix: Classification of global solutions in plane

We are going to classify the homogeneous solutions of degree κ of [\(3\)](#page-1-3) in plane.

Proposition 6.1. Let $n = 2$. If **u** is a homogeneous solutions of degree κ of [\(3\)](#page-1-3) such that $\{x : |\mathbf{u}(x)| = |\nabla \mathbf{u}(x)| = 0\} \neq \{0\}$, then there exist a unit vector $\nu \in \mathbb{R}^2$ and vectors $\mathbf{e}_{+}, \mathbf{e}_{-} \in \mathbb{R}^{m}$ such that $|\mathbf{e}_{\pm}| = 0$ or 1, and

$$
\mathbf{u}(x) = \alpha \max(x \cdot \nu, 0)^{\kappa} \mathbf{e}_+ + \alpha \max(-x \cdot \nu, 0)^{\kappa} \mathbf{e}_-.
$$

Proof. Suppose $\mathbf{u}(x) = r^{\kappa} \Phi(\theta)$, and rewrite the equation in polar coordinates. Then

$$
\Phi'' + \kappa^2 \Phi = f(\Phi), \quad \Phi(0) = \Phi(2\pi).
$$

Use polar coordinates for \mathbb{R}^m and write $\Phi(\theta) = \rho(\theta)\hat{\Phi}(\theta)$, where $\hat{\Phi}$ is a unit vector in \mathbb{R}^m . According to the assumption there is θ_0 such that $\Phi(\theta_0) = \Phi'(\theta_0) = 0$. By a translation we may consider $(0, a)$ to be the maximal interval in which $|\Phi(\theta)| \neq 0$ and $\Phi(0) = \Phi'(0) = 0$. Then ρ and $\hat{\Phi}$ are smooth in $(0, a)$ and satisfy

(26)
$$
\rho'' \hat{\Phi} + 2\rho' \hat{\Phi}' + \rho \hat{\Phi}'' + \kappa^2 \rho \hat{\Phi} = \rho^q \hat{\Phi}.
$$

Now, using this fact that $\hat{\Phi} \cdot \hat{\Phi}' = 0$, and multiplying [\(26\)](#page-20-2) in $\hat{\Phi}'$, we obtain

$$
2\rho'|\hat{\Phi}'|^2 + \rho \hat{\Phi}'' \cdot \hat{\Phi}' = 0.
$$

Thus

$$
\frac{d}{d\theta}(\rho^4|\hat{\Phi}'|^2) = 0,
$$

and $\rho^4|\hat{\Phi}'|^2$ is a constant function in interval $(0, a)$. On the other hand, $\rho(0)$ = $|\Phi(0)| = 0$, and

$$
\rho \hat{\Phi}' = \Phi' - (\hat{\Phi} \cdot \Phi') \hat{\Phi}
$$

is bounded in $(0, a)$. Then $\lim_{\theta \to 0} \rho^4 |\hat{\Phi}'|^2 = 0$, hence $\rho^4 |\hat{\Phi}'|^2 \equiv 0$ for $\theta \in (0, a)$ and so $\hat{\Phi}$ is a constant vector in this interval since $\rho > 0$. (Note that $\hat{\Phi}$ is smooth as long as $\Phi \neq 0.$)

Therefore ρ must satisfy the following equation:

$$
\rho'' + \kappa^2 \rho = \rho^q \quad \text{in } (0, a).
$$

Since $\rho(0) = \rho'(0) = 0$, according to Proposition 3.2 in [\[4\]](#page-21-2), $a = \pi$ and $\rho(\theta) =$ $\alpha \sin^{k}(\theta)$. Then $\Phi(\pi) = 0$. Although ρ' is not necessarily continuous at $\theta = \pi$, Φ is smooth and

$$
\Phi'(\pi) = \lim_{\theta \to \pi^-} \Phi'(\theta) = \lim_{\theta \to \pi^-} \rho'(\theta) \hat{\Phi} = 0.
$$

Now we can repeat the above argument for another maximal interval $(b, c) \subseteq$ $(\pi, 2\pi)$, such that $\Phi(b) = \Phi'(b) = 0$. We thus conclude that either $\rho = 0$ in $(\pi, 2\pi)$ or $(b, c) = (\pi, 2\pi)$ and $\rho(\theta) = \alpha \sin^{k}(\theta - \pi)$. Therefore there are two unit vectors $\hat{\Phi}_+$ and $\hat{\Phi}_-$ such that

$$
\mathbf{u}(x) = \alpha (x_1^+)^{\kappa} \hat{\Phi}_+ + \alpha (x_1^-)^{\kappa} \hat{\Phi}_-.
$$

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